

Conditional estimates in three-dimensional chemotaxis-Stokes systems and application to a Keller-Segel-fluid model accounting for gradient-dependent flux limitation

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Abstract

This manuscript is concerned with the Keller-Segel-Stokes system

$$\begin{cases} n_t + u \cdot \nabla n &= \Delta n - \nabla \cdot (nF(|\nabla c|^2)\nabla c), \\ c_t + u \cdot \nabla c &= \Delta c - c + n, \\ u_t &= \Delta u + \nabla P + n\nabla\Phi, \quad \nabla \cdot u = 0, \end{cases} \quad (\star)$$

under no-flux/no-flux/Dirichlet boundary conditions in smoothly bounded three-dimensional domains, with given suitably regular functions F and Φ . Here in accordance with recent developments in the literature on refined modeling of chemotactic migration, the introduction of suitably decaying F is supposed to adequately account for saturation mechanisms that limit cross-diffusive fluxes near regions of large signal gradients. In the context of such nonlinearities which suitably generalize the prototype given by $F(\xi) = K_F(1 + \xi)^{-\frac{\alpha}{2}}$, $\xi \geq 0$, with $K_F > 0$, known results addressing a fluid-free parabolic-elliptic simplification of (\star) have identified the value $\alpha_c = \frac{1}{2}$ as critical with regard to the occurrence of blow-up in the sense that some exploding solutions can be found when $\alpha < \frac{1}{2}$, whereas all suitably regular initial data give rise to global bounded solutions when $\alpha > \frac{1}{2}$.

The intention of the present study consists in making sure that the latter feature of blow-up prevention by suitably strong flux limitation persists also in the more complex framework of the fully coupled chemotaxis-fluid system (\star) . To achieve this, as a secondary objective of possibly independent interest the manuscript separately establishes some conditional bounds for corresponding fluid fields and taxis gradients in a fairly general setting that particularly includes the subsystem of (\star) concerned with the evolution of (c, u, P) . These estimates relate respective regularity features to certain integrability properties of associated forcing terms, as in the context of (\star) essentially represented by the quantity n .

The application of this tool to the specific problem under consideration thereafter facilitates the derivation of a result on global existence of bounded classical solutions to (\star) for widely arbitrary initial data actually within the entire range $\alpha > \frac{1}{2}$, and by means of an argument which appears to be significantly condensed when compared to reasonings pursued in previous works concerned with related problems.

Key words: chemotaxis; Stokes; flux limitation

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1 Introduction

Flux-limited Keller-Segel systems in liquid environments. Due to their relevance in refined models for chemotactic movement ([2], [32], [6], [7]), dependencies of cell migration rates on gradients of the unknown quantities have received increasing interest in the recent literature on mathematical analysis of taxis-type cross-diffusion systems. Exemplary studies are concerned with global solvability in models involving p -Laplace type cell diffusion operators ([5], [25], [45], [33], [34], [26]), quite precisely identify critical parameters for the occurrence of blow-up in systems simultaneously accounting for flux-limited cross-diffusion and diffusion mechanisms paralleling those in relativistic heat equations ([3], [4], [10], [30]), or also discuss wave-like solution behavior and pattern-supporting instability features in various particular modeling contexts based on flux-limiting mechanisms ([1], [8], [32]).

With regard to issues related to singular structure formation, among the model classes most comprehensively understood seems to be the family of parabolic-elliptic systems given by

$$\begin{cases} n_t = \Delta n - \nabla \cdot (nF(|\nabla c|^2)\nabla c), \\ 0 = \Delta c - c + n, \end{cases} \quad (1.1)$$

where the given function F describes gradient-dependent limitation of cross-diffusive fluxes by suitably generalizing the prototype

$$F(\xi) = K_F(1 + \xi)^{-\frac{\alpha}{2}}, \quad \xi \geq 0, \quad (1.2)$$

with $K_F > 0$ and $\alpha > 0$. Indeed, with respect to the fundamental question how far even despite the presence of such saturation effects, concentration phenomena in the sense of spontaneous emergence of locally unbounded densities may occur in (1.1), when considered under homogeneous Neumann boundary conditions in smoothly bounded N -dimensional domains with $N \geq 2$, the number $\alpha_c(N) := \frac{N-2}{N-1}$ has been found to play the role of a critical exponent: Whenever F is suitably smooth and such that

$$|F(\xi)| \leq K_F \cdot (\xi + 1)^{-\frac{\alpha}{2}} \quad \text{for all } \xi \geq 0 \quad (1.3)$$

with some $K_F > 0$ and $\alpha > \alpha_c$, an associated initial-boundary value problem admits globally bounded classical solutions for widely arbitrary initial data, while if $N \geq 3$, Ω is a ball and

$$F(\xi) \geq k_F \cdot (\xi + 1)^{-\frac{\alpha}{2}} \quad \text{for all } \xi \geq 0 \quad (1.4)$$

with some $k_F > 0$ and $\alpha < \alpha_c$, then some radially symmetric classical solutions exist which blow up at some finite time $T > 0$ in the sense that $\limsup_{t \nearrow T} \|n(\cdot, t)\|_{L^\infty(\Omega)} = \infty$ ([43]; cf. also [31] and [44]).

Having this fairly complete knowledge for (1.1) at hand, in this study we intend to accomplish an apparently natural next step toward an understanding of models for tactic movement in realistic situations: Namely, we plan to examine to which extent the above blow-up related dichotomy may be affected if buoyancy-induced interaction of a considered population with its liquid environment is additionally accounted for. In fact, experimental evidence indicates a substantial relevance of chemotaxis-fluid couplings in several contexts ([12], [13], [29], [35], [36]), and both numerical ([36], [28]) and some rigorous analytical precedents ([22], [23], [24], [20], [14]) witness significant quantitative effects of fluid interaction in various specific situations, mostly involving a given and hence passive fluid.

To take aim at this question in a fairly general three-dimensional setting compatible with the seminal modeling approach in [36], we shall below consider a fully parabolic counterpart of (1.1), coupled to the incompressible Stokes equations, in the framework of the initial-boundary value problem

$$\begin{cases} n_t + u \cdot \nabla n &= \Delta n - \nabla \cdot (nF(|\nabla c|^2)\nabla c), & x \in \Omega, t > 0, \\ c_t + u \cdot \nabla c &= \Delta c - c + n, & x \in \Omega, t > 0, \\ u_t &= \Delta u + \nabla P + n\nabla\Phi, \quad \nabla \cdot u = 0, & x \in \Omega, t > 0, \\ \frac{\partial n}{\partial \nu} = \frac{\partial c}{\partial \nu} = 0, \quad u = 0, & & x \in \partial\Omega, t > 0, \\ n(x, 0) = n_0(x), \quad c(x, 0) = c_0(x), \quad u(x, 0) = u_0(x), & & x \in \Omega, \end{cases} \quad (1.5)$$

where $\Omega \subset \mathbb{R}^3$ is a bounded domain with smooth boundary, where $\Phi \in W^{2,\infty}(\Omega)$ represents a given gravitational potential, and where F appropriately generalizes the choice in (1.2). Since any exploding solution of the corresponding fluid-free chemotaxis-only system trivially extends to an accordingly singular solution of (1.5) upon letting $\Phi \equiv \text{const.}$ and $u_0 \equiv 0$ therein, and since in view of the above criticality feature of $\alpha_c(3) = \frac{1}{2}$ we thus expect blow-up to occur in (1.5) whenever F satisfies (1.4) with some $k_F > 0$ and $\alpha < \frac{1}{2}$, the major part of our ambition seems to consist in deciding whether or not the assumption (1.3) with $K_F > 0$ and arbitrary $\alpha > \frac{1}{2}$ remains sufficient to ensure global smooth solvability and boundedness also in the context of the full problem (1.5).

Conditional L^∞ estimates for signal gradients in chemotaxis-Stokes systems. Main results

I. In order to develop an approach which is not only appropriate to address this topic, but which is suitably broad so as to bear some potential for meaningful applications in different frameworks beyond the particular present one, prior to our analysis of the specific system (1.5) we shall take up the general problem of efficiently controlling signal gradients in contexts of chemotaxis processes coupled to Stokes fluids. This problem has formed a natural core of several precedent studies concerned with models related to (1.5), and a considerable variety of respective ad-hoc approaches, based either on the detection of certain entropy-like features, or on maximal Sobolev regularity results, or on heat semigroup estimates, can be found in the literature (cf. [27], [37], [38], [9], [41] for examples).

An essential characteristic of our subsequent analysis will now be made up by the intention to propose the use of certain conditional L^∞ estimates for taxis gradients as a possible alternative to previous strategies. The derivation of such bounds in suitably general settings will form a fairly extensive first part of this manuscript, but as a considerable benefit of the related efforts, inter alia indicating their potential value beyond the context of (1.5), we will thereafter be able to conveniently accomplish our original objective by means of quite a compressed argument.

In order to make this more precise, in a bounded domain $\Omega \subset \mathbb{R}^3$ and with given $T \in (0, \infty]$, $m > 0, K > 0$ and $\vartheta \in (\frac{3}{4}, 1)$, let us consider solutions (z, v, Q) , with

$$\begin{cases} z \in \bigcap_{q>3} C^0([0, T]; W^{1,q}(\Omega)) \cap C^{2,1}(\bar{\Omega} \times (0, T)), \\ v \in C^0([0, T]; D(A^\vartheta)) \cap C^{2,1}(\bar{\Omega} \times (0, T); \mathbb{R}^3) \quad \text{and} \\ Q \in C^{1,0}(\Omega \times (0, T)), \end{cases} \quad (1.6)$$

of the problem

$$\begin{cases} z_t + v \cdot \nabla z = \Delta z - z + f, & x \in \Omega, t \in (0, T), \\ v_t = \Delta v + \nabla Q + f \nabla \Phi, \quad \nabla \cdot v = 0, & x \in \Omega, t \in (0, T), \\ \frac{\partial z}{\partial \nu} = 0, \quad v = 0, & x \in \partial\Omega, t \in (0, T), \\ z(x, 0) = z_0(x), \quad v(x, 0) = v_0(x), & x \in \Omega, \end{cases} \quad (1.7)$$

where z_0, v_0 and f are given functions which are such that

$$\begin{cases} z_0 \in W^{1,\infty}(\Omega) \text{ is nonnegative and } v_0 \in D(A^\vartheta) & \text{with} \\ \|z_0\|_{W^{1,\infty}(\Omega)} \leq K \text{ and } \|A^\vartheta v_0\|_{L^2(\Omega)} \leq K, \end{cases} \quad (1.8)$$

and that

$$\begin{cases} f \in C^0(\overline{\Omega} \times [0, T]) \text{ is nonnegative with} \\ \int_{\Omega} f(\cdot, t) \leq m \text{ for all } t \in (0, T). \end{cases} \quad (1.9)$$

Here, as throughout the sequel, we let $A = -\mathcal{P}\Delta$ and $(A^\vartheta)_{\vartheta>0}$ denote the realization of the Stokes operator in $L^2(\Omega; \mathbb{R}^3)$, with domain given by $D(A) := W^{2,2}(\Omega; \mathbb{R}^3) \cap W_{0,\sigma}^{1,2}(\Omega)$, and the family of its corresponding fractional powers, respectively, where \mathcal{P} represents the Helmholtz projection on $L^2(\Omega; \mathbb{R}^3)$, and where $W_{0,\sigma}^{1,2}(\Omega) := W_0^{1,2}(\Omega; \mathbb{R}^3) \cap L_\sigma^2(\Omega)$ with $L_\sigma^2(\Omega) := \{\varphi \in L^2(\Omega; \mathbb{R}^3) \mid \nabla \cdot \varphi = 0\}$.

Within this framework, by means of quite a direct application of standard regularity estimates for the Stokes semigroup we shall firstly obtain the following statement on regularity of the corresponding fluid field, conditional in that some temporally independent L^p bounds for the forcing term f in (1.7) are involved.

Proposition 1.1 *Let $\Omega \subset \mathbb{R}^3$ be a bounded domain with smooth boundary, and suppose that $\vartheta \in (\frac{3}{4}, 1)$, $m > 0, K > 0, p \geq 2, \theta \in (\frac{1}{4}, \vartheta]$ and $\eta > 0$. Then there exists $C(m, K, p, \theta, \eta) > 0$ such that any f, z, v and Q fulfilling (1.9), (1.6) and (1.7), with some $T \in (0, \infty]$ and some (z_0, v_0) complying with (1.8), satisfies*

$$\|A^\theta v(\cdot, t)\|_{L^2(\Omega)} \leq C(m, K, p, \theta, \eta) \cdot \left\{ 1 + \sup_{s \in (0, t)} \|f(\cdot, s)\|_{L^p(\Omega)} \right\}^{\frac{p}{p-1} \cdot (\frac{4\theta-1}{6} + \eta)} \quad \text{for all } t \in (0, T). \quad (1.10)$$

On the basis thereof and of appropriate regularization features of Neumann heat semigroups, we shall secondly derive uniform estimates for the first solution components in (1.7) and their gradients, subject to a conditionality property similar to that expressed in (1.10):

Theorem 1.2 *Suppose that $\Omega \subset \mathbb{R}^3$ is a bounded domain with smooth boundary, and let $\vartheta \in (\frac{3}{4}, 1)$, $m > 0, K > 0, p > 3$ and $\eta > 0$. Then one can find $C(m, K, p, \eta) > 0$ with the property that given any $T \in (0, \infty]$ as well as functions f, z, v and Q which satisfy (1.9), (1.6) and (1.7) with some (z_0, v_0) fulfilling (1.8), we have*

$$\|z(\cdot, t)\|_{W^{1,\infty}(\Omega)} \leq C(m, K, p, \eta) \cdot \left\{ 1 + \sup_{s \in (0, t)} \|f(\cdot, s)\|_{L^p(\Omega)} \right\}^{\frac{p}{p-1} \cdot (\frac{2}{3} + \eta)} \quad \text{for all } t \in (0, T). \quad (1.11)$$

Approaching criticality in boundedness results for (1.5). Main results II. With the above preparations at hand, in Section 3 we can thereafter return to our original question. In fact, Proposition 1.1 and, especially, Theorem 1.2 will enable us to make sure by means of a fairly condensed argument that global existence and boundedness in (1.5) can indeed be achieved under the assumption (1.3) within the entire range of α known as the regime essentially maximal in this regard for (1.5):

Theorem 1.3 *Let $\Omega \subset \mathbb{R}^3$ be a bounded domain with smooth boundary, and let $\Phi \in W^{2,\infty}(\Omega)$ and $F \in C^2([0, \infty))$ be such that*

$$|F(\xi)| \leq K_F \cdot (1 + \xi)^{-\frac{\alpha}{2}} \quad \text{for all } \xi \geq 0 \quad (1.12)$$

with some $K_F > 0$ and

$$\alpha > \frac{1}{2}.$$

Then for any

$$\begin{cases} n_0 \in C^0(\overline{\Omega}) \text{ such that } n_0 \geq 0, \\ c_0 \in W^{1,\infty}(\Omega) \text{ such that } c_0 \geq 0 \quad \text{and} \\ u_0 \in \bigcup_{\vartheta \in (\frac{3}{4}, 1)} D(A^\vartheta), \end{cases} \quad (1.13)$$

there exist uniquely determined

$$\begin{cases} n \in C^0(\overline{\Omega} \times [0, \infty)) \cap C^{2,1}(\overline{\Omega} \times (0, \infty)), \\ c \in \bigcap_{q>3} C^0([0, \infty); W^{1,q}(\Omega)) \cap C^{2,1}(\overline{\Omega} \times (0, \infty)) \quad \text{and} \\ u \in \bigcup_{\vartheta \in (\frac{3}{4}, 1)} C^0([0, \infty); D(A^\vartheta)) \cap C^{2,1}(\overline{\Omega} \times (0, \infty); \mathbb{R}^3) \end{cases}$$

such that $n \geq 0$ and $c \geq 0$ in $\Omega \times (0, \infty)$, and that with some $P \in C^{1,0}(\Omega \times (0, \infty))$, the collection (n, c, u, P) solves (1.5) in the classical sense in $\Omega \times (0, \infty)$. Moreover, this solution is bounded in the sense that with some $\vartheta \in (\frac{3}{4}, 1)$ and $C > 0$,

$$\|n(\cdot, t)\|_{L^\infty(\Omega)} + \|c(\cdot, t)\|_{W^{1,\infty}(\Omega)} + \|A^\vartheta u(\cdot, t)\|_{L^2(\Omega)} \leq C \quad \text{for all } t > 0. \quad (1.14)$$

Before going into details, let us once again emphasize that applications of the above general estimates, and especially that from Theorem 1.2, to much a larger variety of chemotaxis-Stokes systems beyond the particular context of (1.5) are conceivable. Although in order to keep this manuscript reasonably focused we do not discuss this here in detail, we at least may announce that one further example concerned with a quasilinear Keller-Segel system coupled to the Stokes equations will be addressed in [42].

2 Conditional bounds. Proofs of Proposition 1.1 and Theorem 1.2

2.1 Fluid regularity

Let us first employ known smoothing properties of the Stokes semigroup to see that in the considered context of solutions (z, v, Q) of (1.7), regularity properties of the respective fluid can be related to the

respective forcing term through the expressions

$$M_p(t) := 1 + \sup_{s \in (0,t)} \|f(\cdot, s)\|_{L^p(\Omega)}, \quad t \in (0, T), \quad (2.1)$$

well-defined due to (1.9) actually for all $p \geq 1$, in the quantitative manner specified in Proposition 1.1:

PROOF of Proposition 1.1. Since $\theta < 1$, we may assume that $\eta < \frac{2(1-\theta)}{3}$, which ensures that $7 - 4\theta - 6\eta > 3$, whence observing that $7 - 4\theta < 6$ due to the restriction $\theta > \frac{1}{4}$, we obtain that

$$1 < \frac{6}{7 - 4\theta} < \frac{6}{7 - 4\theta - 6\eta} < 2.$$

As, on the other hand, we are presupposing that $p \geq 2$, if we choose any

$$\lambda = \lambda(p, \theta, \eta) \in \left(\frac{6}{7 - 4\theta}, \frac{6}{7 - 4\theta - 6\eta} \right], \quad (2.2)$$

then necessarily $1 < \lambda < 2 \leq p$, so that by continuity of the Helmholtz projection on $L^\lambda(\Omega; \mathbb{R}^3)$ ([17]), we can find $C_1 = C_1(p, \theta, \eta) > 0$ such that abbreviating $C_2 := \|\nabla\Phi\|_{L^\infty(\Omega)}$ and $a := \frac{p(\lambda-1)}{(p-1)\lambda}$, due to the Hölder inequality we have

$$\begin{aligned} \|\mathcal{P}[\varphi\nabla\Phi]\|_{L^\lambda(\Omega)} &\leq C_1\|\varphi\nabla\Phi\|_{L^\lambda(\Omega)} \\ &\leq C_1C_2\|\varphi\|_{L^\lambda(\Omega)} \\ &\leq C_1C_2\|\varphi\|_{L^p(\Omega)}^a\|\varphi\|_{L^1(\Omega)}^{1-a} \quad \text{for all } \varphi \in C^0(\bar{\Omega}). \end{aligned} \quad (2.3)$$

We next use that $\theta \leq \vartheta$ and invoke known regularization features of the Stokes semigroup $(e^{-tA})_{t \geq 0}$ ([19, p.201]) to infer that with some $C_3 = C_3(\theta) > 0$, $C_4 = C_4(p, \theta, \eta) > 0$ and $\mu > 0$,

$$\|A^\theta e^{-tA}\varphi\|_{L^2(\Omega)} \leq C_3\|A^\vartheta\varphi\|_{L^2(\Omega)} \quad \text{for all } t \geq 0 \text{ and } \varphi \in D(A^\vartheta) \quad (2.4)$$

as well as

$$\|A^\theta e^{-tA}\varphi\|_{L^2(\Omega)} \leq C_4 t^{-\theta - \frac{3}{2}(\frac{1}{\lambda} - \frac{1}{2})} e^{-\mu t} \|\varphi\|_{L^\lambda(\Omega)} \quad \text{for all } t > 0 \text{ and } \varphi \in C^0(\bar{\Omega}; \mathbb{R}^3). \quad (2.5)$$

We now let $m > 0$ and $K > 0$ be given, and suppose that f, z_0, v_0, v, z and Q satisfy (1.6), (1.7), (1.9) and (1.8) with some $T \in (0, \infty]$. Then taking $(M_p(t))_{t \in (0, T)}$ as accordingly introduced in (2.1), from (2.3) we obtain that

$$\|\mathcal{P}[f(\cdot, s)\nabla\Phi]\|_{L^\lambda(\Omega)} \leq C_1C_2m^{1-a}M_p^a(t) \quad \text{for all } t \in (0, T) \text{ and any } s \in (0, t),$$

whence a combination of (2.4) with (2.5) shows that

$$\begin{aligned} \|A^\theta v(\cdot, t)\|_{L^2(\Omega)} &= \left\| A^\theta e^{-tA}v_0 + \int_0^t A^\theta e^{-(t-s)A}\mathcal{P}[f(\cdot, s)\nabla\Phi]ds \right\|_{L^2(\Omega)} \\ &\leq C_3\|A^\vartheta v_0\|_{L^2(\Omega)} + C_4 \int_0^t (t-s)^{-\theta - \frac{3}{2}(\frac{1}{\lambda} - \frac{1}{2})} e^{-\mu(t-s)} \|\mathcal{P}[f(\cdot, s)\nabla\Phi]\|_{L^\lambda(\Omega)} ds \\ &\leq C_3K + C_1C_2C_4C_5m^{1-a}M_p^a(t) \quad \text{for all } t \in (0, T) \end{aligned} \quad (2.6)$$

thanks to (1.8), where $C_5 = C_5(p, \theta, \eta) := \int_0^\infty \sigma^{-\theta - \frac{3}{2}(\frac{1}{\lambda} - \frac{1}{2})} e^{-\mu\sigma} d\sigma$ is finite due to the lower bound for λ implied by (2.2), which namely warrants that

$$\theta + \frac{3}{2} \cdot \left(\frac{1}{\lambda} - \frac{1}{2} \right) < \theta + \frac{3}{2} \cdot \left(\frac{7 - 4\theta}{6} - \frac{1}{2} \right) = 1.$$

Since, conversely, the upper bound for λ in (2.2) guarantees that

$$a = \frac{p}{p-1} \cdot \left(1 - \frac{1}{\lambda} \right) \leq \frac{p}{p-1} \cdot \left(1 - \frac{7 - 4\theta - 6\eta}{6} \right) = \frac{p}{p-1} \cdot \left(\frac{4\theta - 1}{6} + \eta \right),$$

and that thus the inequality $M_p \geq 1$ ensures that $M_p^a(t) \leq M_p^{\frac{p}{p-1} \cdot (\frac{4\theta-1}{6} + \eta)}(t)$ for all $t \in (0, T)$, we conclude that (1.10) is a consequence of (2.6). \square

In order to appropriately prepare an application of the latter in the course of our derivation of Theorem 1.2, we recall known embedding properties enjoyed by the domains of the operators A^θ appearing in (1.10) to turn the latter into integral bounds for the fluid velocity itself:

Corollary 2.1 *Let $\vartheta \in (\frac{3}{4}, 1)$, and let $m > 0, K > 0, p \geq 2, r > 3$ and $\eta > 0$. Then there exists $C(m, K, p, r, \eta) > 0$ such that if (1.9) and (1.8) as well as (1.6) and (1.7) hold with some $T \in (0, \infty]$, it follows that*

$$\|v(\cdot, t)\|_{L^r(\Omega)} \leq C(m, K, p, r, \eta) \cdot \left\{ 1 + \sup_{s \in (0, t)} \|f(\cdot, s)\|_{L^p(\Omega)} \right\}^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta)} \quad \text{for all } t \in (0, T). \quad (2.7)$$

PROOF. Since $4\vartheta - 3 \geq 0$ and thus $\frac{(4\vartheta-3)r+6}{5r} > 0$, we may restrict ourselves to considering the case when $\eta \leq \frac{(4\vartheta-3)r+6}{5r}$, in which

$$\theta = \theta(r, \eta) := \frac{3r - 6 + 5r\eta}{4r} \quad (2.8)$$

satisfies

$$\theta \leq \frac{3r - 6 + 5r \cdot \frac{(4\vartheta-3)r+6}{5r}}{4r} = \vartheta.$$

As (2.8) furthermore clearly ensures that

$$\theta > \frac{3r - 6}{4r}, \quad (2.9)$$

and that thus

$$\theta > \frac{3 - \frac{6}{r}}{4} > \frac{1}{4}$$

due to our assumption that $r > 3$, we may employ Proposition 1.1 to find $C_1 = C_1(m, K, p, r, \eta)$ with the property that if $T \in (0, \infty]$ as well as (1.9), (1.8), (1.6) and (1.7) hold, then with $(M_p(t))_{t \in (0, T)}$ taken from (2.1) we have

$$\|A^\theta v(\cdot, t)\|_{L^2(\Omega)} \leq C_1 M_p^{\frac{p}{p-1} \cdot \frac{4\theta-1+\eta}{6}}(t) \quad \text{for all } t \in (0, T). \quad (2.10)$$

Apart from that, (2.9) implies that

$$2\theta - \frac{3}{2} + \frac{3}{r} > 2 \cdot \frac{3r-6}{4r} - \frac{3}{2} + \frac{3}{r} = 0,$$

from which it follows that $D(A^\theta)$ is continuously embedded into $L^r(\Omega; \mathbb{R}^3)$ ([18], [21]), and that hence we can pick $C_2 = C_2(r, \eta) > 0$ such that

$$\|\varphi\|_{L^r(\Omega)} \leq C_2 \|A^\theta \varphi\|_{L^2(\Omega)} \quad \text{for all } \varphi \in D(A^\theta). \quad (2.11)$$

It now only remains to observe that if we assume (1.9), (1.8), (1.6) and (1.7) to be satisfied with some $T \in (0, \infty]$, then a combination of (2.11) with (2.10) shows that with M_p as accordingly defined through (2.1) we have

$$\|v(\cdot, t)\|_{L^r(\Omega)} \leq C_1 C_2 M_p^{\frac{p}{p-1} \cdot \frac{4\theta-1+\eta}{6}}(t) \quad \text{for all } t \in (0, T),$$

and that

$$\frac{4\theta - 1 + \eta}{6} = \frac{4 \cdot \frac{3r-6+5r\eta}{4r} - 1 + \eta}{6} = \frac{2r - 6 + 6r\eta}{6r} = \frac{r-3}{3r} + \eta$$

by (2.8). □

2.2 Uniform estimates for taxis gradients

The first step toward our analysis concerned with the first equation in (1.7) acts at levels which, through relying on a standard zero-order testing procedure only, do not explicitly rely on any property of the considered fluid field beyond its mere solenoidality. Although the estimates for z thereby obtained involve topologies which are yet quite far from those to be finally addressed, they will form an essential basic information in the course of an interpolation argument performed in Lemma 2.3 below.

Lemma 2.2 *Let $\vartheta \in (\frac{3}{4}, 1)$. Then for any $m > 0, K > 0, p \geq 2$ and $q \geq 2$ there exists $C(m, K, p, q) > 0$ such that if (1.9), (1.8), (1.6) and (1.7) hold with some $T \in (0, \infty]$, then*

$$\|z(\cdot, t)\|_{L^q(\Omega)} \leq C(m, K, p, q) \cdot \left\{ 1 + \sup_{s \in (0, t)} \|f(\cdot, s)\|_{L^p(\Omega)} \right\}^{\frac{p}{p-1} \cdot \frac{q-1}{3q}} \quad \text{for all } t \in (0, T). \quad (2.12)$$

PROOF. By continuity of the embedding $W^{1,2}(\Omega) \hookrightarrow L^6(\Omega)$, we can fix $C_1 > 0$ such that

$$\|\varphi\|_{L^6(\Omega)}^2 \leq C_1 \int_{\Omega} |\nabla \varphi|^2 + C_1 \int_{\Omega} \varphi^2 \quad \text{for all } \varphi \in W^{1,2}(\Omega). \quad (2.13)$$

Therefore, assuming (1.9), (1.8), (1.6) and (1.7) with some $T \in (0, \infty]$, in the identity

$$\frac{1}{q} \frac{d}{dt} \int_{\Omega} z^q + \frac{4(q-1)}{q^2} \int_{\Omega} |\nabla z^{\frac{q}{2}}|^2 + \int_{\Omega} z^q = \int_{\Omega} z^{q-1} f, \quad t \in (0, T), \quad (2.14)$$

as obtained upon testing the first equation in (1.7) by z^{q-1} due to the solenoidality of v , we may combine the Hölder inequality with (2.13) and Young's inequality to estimate

$$\begin{aligned}
\int_{\Omega} z^{q-1} f &= \int_{\Omega} (z^{\frac{q}{2}})^{\frac{2(q-1)}{q}} f \\
&\leq \|z^{\frac{q}{2}}\|_{L^6(\Omega)}^{\frac{2(q-1)}{q}} \|f\|_{L^{\frac{3q}{2q+1}}(\Omega)} \\
&\leq C_1^{\frac{q-1}{q}} \cdot \left\{ \int_{\Omega} |\nabla z^{\frac{q}{2}}|^2 + \int_{\Omega} z^q \right\}^{\frac{q-1}{q}} \cdot \|f\|_{L^{\frac{3q}{2q+1}}(\Omega)} \\
&= \left\{ \frac{1}{q} \cdot \left\{ \int_{\Omega} |\nabla z^{\frac{q}{2}}|^2 + \int_{\Omega} z^q \right\}^{\frac{q-1}{q}} \right\} \cdot q^{\frac{q-1}{q}} C_1^{\frac{q-1}{q}} \|f\|_{L^{\frac{3q}{2q+1}}(\Omega)} \\
&\leq \frac{q-1}{q} \cdot \left\{ \frac{1}{q} \cdot \left\{ \int_{\Omega} |\nabla z^{\frac{q}{2}}|^2 + \int_{\Omega} z^q \right\} \right\} + \frac{1}{q} \cdot \left\{ q^{\frac{q-1}{q}} C_1^{\frac{q-1}{q}} \|f\|_{L^{\frac{3q}{2q+1}}(\Omega)} \right\}^q \\
&= \frac{q-1}{q^2} \int_{\Omega} |\nabla z^{\frac{q}{2}}|^2 + \frac{q-1}{q^2} \int_{\Omega} z^q + q^{q-2} C_1^{q-1} \|f\|_{L^{\frac{3q}{2q+1}}(\Omega)}^q \quad \text{for all } t \in (0, T). \tag{2.15}
\end{aligned}$$

Here since $1 < \frac{3q}{2q+1} < \frac{3}{2} < p$, we may once again rely on the Hölder inequality to see that thanks to (1.9), with $(M_p(t))_{t \in (0, T)}$ as in (2.1) we have

$$\begin{aligned}
\|f\|_{L^{\frac{3q}{2q+1}}(\Omega)}^q &\leq \|f\|_{L^p(\Omega)}^{\frac{p(q-1)}{3(p-1)}} \|f\|_{L^1(\Omega)}^{\frac{2pq+p-3q}{3(p-1)}} \\
&\leq m^{\frac{2pq+p-3q}{3(p-1)}} M_p^{\frac{p(q-1)}{3(p-1)}}(t) \quad \text{for all } t \in (0, T).
\end{aligned}$$

As

$$\frac{4(q-1)}{q^2} \geq \frac{q-1}{q^2} \quad \text{and} \quad 1 - \frac{q-1}{q^2} \geq 1 - \frac{1}{q} \geq \frac{1}{2} \geq \frac{1}{q}$$

due to the fact that $q \geq 2$, from (2.14) and (2.15) we therefore obtain that for any choice of $t_0 \in (0, T)$,

$$\frac{d}{dt} \int_{\Omega} z^q + \int_{\Omega} z^q \leq q^{q-1} C_1^{q-1} m^{\frac{2pq+p-3q}{3(p-1)}} M_p^{\frac{p(q-1)}{3(p-1)}}(t_0) \quad \text{for all } t \in (0, t_0),$$

and that thus, according to an ODE comparison argument,

$$\int_{\Omega} z^q \leq \max \left\{ \int_{\Omega} z_0^q, q^{q-1} C_1^{q-1} m^{\frac{2pq+p-3q}{3(p-1)}} M_p^{\frac{p(q-1)}{3(p-1)}}(t_0) \right\} \quad \text{for all } t \in (0, t_0].$$

When evaluated at $t = t_0$, in view of (1.8) this readily yields the claim. \square

Indeed, the latter can be used to derive the following preliminary information on control of spatial $W^{1, \infty}$ norms of z , as addressed in Theorem 1.2. In its formulation and throughout the remainder of this section, we let $B = B_q$ denote the sectorial realization of $-\Delta + 1$ under homogeneous Neumann boundary conditions in $\bigcap_{q>1} L^q(\Omega)$, and let $(B^\beta)_{\beta>0}$ represent the associated family of positive

fractional powers $B^\beta = B_q^\beta$. Then the respective domains $D(B_q^\beta)$ are continuously embedded into $W^{1,\infty}(\Omega)$ whenever $2\beta - \frac{3}{q} > 1$, whence for such parameters the quantities under consideration can be estimated against expressions of the form $\|B^\beta z\|_{L^q(\Omega)}$, together with the Lebesgue norms from Lemma 2.2, through interpolation. This is substantiated in the following statement which circumvents problems potentially resulting from possibly insufficient regularity properties of z_0 , as merely assumed here to belong to $W^{1,\infty}(\Omega)$ but not necessarily to $D(B_q^\beta)$ for any $\beta > \frac{1}{2}$, by subtracting a suitable correction.

Lemma 2.3 *Let $\vartheta \in (\frac{3}{4}, 1)$, and let $m > 0, K > 0, p \geq 2, q \geq 2$ and $\beta \in (\frac{1}{2}, 1)$ be such that*

$$q \cdot (2\beta - 1) > 3. \quad (2.16)$$

Then for all $\eta > 0$ one can find $C(m, K, p, q, \beta, \eta) > 0$ such that whenever (1.9), (1.8), (1.6) and (1.7) are satisfied with some $T \in (0, \infty]$, we have

$$\begin{aligned} \|z(\cdot, t) - e^{-tB} z_0\|_{W^{1,\infty}(\Omega)} &\leq C(m, K, p, q, \beta, \eta) \cdot \left\{ 1 + \sup_{s \in (0, t)} \|f(\cdot, s)\|_{L^p(\Omega)} \right\}^{\frac{p}{p-1} \cdot \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta}} \times \\ &\quad \times \left\{ 1 + \left\| B^\beta \left(z(\cdot, t) - e^{-tB} z_0 \right) \right\|_{L^q(\Omega)} \right\}^{\left(\frac{q+3}{2q\beta} + \eta \right)} \end{aligned} \quad (2.17)$$

for all $t \in (0, T)$.

PROOF. As (2.16) warrants that

$$\frac{q+3}{2q} = \frac{1}{2} + \frac{3}{2q} < \frac{1}{2} + \frac{3}{2} \cdot \frac{2\beta - 1}{3} = \beta,$$

we may assume that η is so small that

$$\gamma = \gamma(q, \eta) := \frac{q+3}{2q} + \beta\eta \quad (2.18)$$

satisfies $\gamma < \beta$. We may therefore draw on a standard interpolation property enjoyed by fractional powers of sectorial operators in quite general settings (see e.g. [15, Theorem 2.14.1]) in choosing $C_1 = C_1(q, \beta, \eta) > 0$ such that

$$\|B^\gamma \varphi\|_{L^q(\Omega)} \leq C_1 \|B^\beta \varphi\|_{L^q(\Omega)}^{\frac{\gamma}{\beta}} \|\varphi\|_{L^q(\Omega)}^{\frac{\beta-\gamma}{\beta}} \quad \text{for all } \varphi \in D(B^\beta). \quad (2.19)$$

Furthermore, the evident fact that $\gamma > \frac{q+3}{2q}$ guarantees that $2\gamma - \frac{3}{q} > 1$, meaning that $D(B^\gamma)$ is continuously embedded into $W^{1,\infty}(\Omega)$ ([21]), and that thus there exists $C_2 = C_2(q, \beta, \eta) > 0$ fulfilling

$$\|\varphi\|_{W^{1,\infty}(\Omega)} \leq C_2 \|B^\gamma \varphi\|_{L^q(\Omega)} \quad \text{for all } \varphi \in D(B^\gamma). \quad (2.20)$$

As a final preparation, we take $C_3 = C_3(q) > 0$ in such a way that

$$\|\varphi\|_{L^q(\Omega)} \leq C_3 \|\varphi\|_{W^{1,\infty}(\Omega)} \quad \text{for all } \varphi \in W^{1,\infty}(\Omega), \quad (2.21)$$

and invoke Lemma 2.2 to see that whenever $m > 0, K > 0$ and $p \geq 2$, we can fix $C_4 = C_4(m, K, p, q) > 0$ with the property that if (1.9), (1.8), (1.6) and (1.7) hold with some $T \in (0, \infty]$, then

$$\|z(\cdot, t)\|_{L^q(\Omega)} \leq C_4 M_p^{\frac{p}{p-1} \cdot \frac{q-1}{3q}}(t) \quad \text{for all } t \in (0, T), \quad (2.22)$$

again with $(M_p(t))_{t \in (0, T)}$ as defined in (2.1).

Thus, given any such T and arbitrary f, z_0, v_0, z, v and Q satisfying (1.9), (1.8), (1.6) and (1.7), abbreviating

$$\widehat{z}(\cdot, t) := z(\cdot, t) - e^{-tB} z_0, \quad t \in (0, T), \quad (2.23)$$

we can firstly rely on (2.22), (2.21) and (1.8) to see that with the correspondingly defined function $M_p \geq 1$, since e^{-tB} is nonexpansive on $L^q(\Omega)$ for all $t > 0$ we have

$$\begin{aligned} \|\widehat{z}(\cdot, t)\|_{L^q(\Omega)} &\leq \|z(\cdot, t)\|_{L^q(\Omega)} + \|e^{-tB} z_0\|_{L^q(\Omega)} \\ &\leq C_4 M_p^{\frac{p}{p-1} \cdot \frac{q-1}{q}}(t) + \|z_0\|_{L^q(\Omega)} \\ &\leq C_4 M_p^{\frac{p}{p-1} \cdot \frac{q-1}{q}}(t) + C_3 K \\ &\leq C_5 M_p^{\frac{p}{p-1} \cdot \frac{q-1}{3q}}(t) \quad \text{for all } t \in (0, T), \end{aligned}$$

where $C_5 = C_5(m, K, p, q) := \max\{C_4, C_3 K\}$. Combining (2.20) with (2.19) hence shows that

$$\begin{aligned} \|\widehat{z}(\cdot, t)\|_{W^{1, \infty}(\Omega)} &\leq C_2 \|B^\gamma \widehat{z}(\cdot, t)\|_{L^q(\Omega)} \\ &\leq C_1 C_2 \|B^\beta \widehat{z}(\cdot, t)\|_{L^q(\Omega)}^{\frac{\gamma}{\beta}} \|\widehat{z}(\cdot, t)\|_{L^q(\Omega)}^{\frac{\beta-\gamma}{\beta}} \\ &\leq C_1 C_2 C_5^{\frac{\beta-\gamma}{\beta}} M_p^{\frac{p}{p-1} \cdot \frac{q-1}{3q} \cdot \frac{\beta-\gamma}{\beta}}(t) \|B^\beta \widehat{z}(\cdot, t)\|_{L^q(\Omega)}^{\frac{\gamma}{\beta}} \quad \text{for all } t \in (0, T), \end{aligned}$$

from which (2.17) results upon observing that

$$\frac{\gamma}{\beta} = \frac{q+3}{2q\beta} + \eta$$

by (2.18), and that

$$\begin{aligned} \frac{p}{p-1} \cdot \frac{q-1}{3q} \cdot \frac{\beta-\gamma}{\beta} &\leq \frac{p}{p-1} \cdot \frac{q-1}{3q} \cdot \frac{\beta - \frac{q+3}{2q}}{\beta} \\ &= \frac{p}{p-1} \cdot \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta} \end{aligned}$$

due to the inequality $\gamma \geq \frac{q+3}{2q}$. □

The core of our analysis in this section is now contained in the following estimate for the higher-order expressions on the right of (2.17) against the quantities in (2.1). This is achieved by means of smoothing estimates for the semigroup $(e^{-tB})_{t \geq 0}$, where the corresponding convective contribution to the first equation in (1.7) will be treated as a lower-order inhomogeneity. Thanks to our preparations

provided by Corollary 2.1 and Lemma 2.3, the resulting influence can appropriately be estimated in terms of the expressions in (2.1), and of the quantities under consideration themselves, namely of

$$N_{q,\beta}(t) := 1 + \sup_{s \in (0,t)} \left\| B^\beta \left(z(\cdot, s) - e^{-sB} z_0 \right) \right\|_{L^q(\Omega)}, \quad t \in (0, T), \quad (2.24)$$

at a conveniently controllable sublinear power (cf. (2.44)). If the fractional power β appearing herein is suitably close to 1, then this indeed yields the following key result on regularity of z .

Lemma 2.4 *Let $\vartheta \in (\frac{3}{4}, 1)$, and suppose that $m > 0, K > 0, p \geq 2, q \geq 2$ and $\beta \in (\frac{1}{2}, 1)$ are such that (2.16) holds, and that*

$$2p(1 - \beta) \geq 3. \quad (2.25)$$

Then for all $\eta > 0$ there exists $C(m, K, p, q, \beta, \eta) > 0$ with the property that if (1.9), (1.8), (1.6) and (1.7) hold with some $T \in (0, \infty]$, then for all $t \in (0, T)$,

$$\left\| B^\beta \left(z(\cdot, t) - e^{-tB} z_0 \right) \right\|_{L^q(\Omega)} \leq C(m, K, p, q, \beta, \eta) \cdot \left\{ 1 + \sup_{s \in (0,t)} \|f(\cdot, s)\|_{L^p(\Omega)} \right\}^{\frac{p}{p-1} \cdot \left(\frac{2q\beta+q-1}{3q} + \eta \right)}. \quad (2.26)$$

PROOF. We prepare our estimation procedure by firstly using that $q \geq \frac{3}{2\beta-1}$ and $p \geq \frac{3}{2(1-\beta)}$ to verify that

$$\lambda = \lambda(p, q, \beta) := \min \left\{ q, \frac{3q}{2q - 2q\beta + 1} \right\} \quad (2.27)$$

satisfies

$$\lambda > \min \left\{ q, \frac{3q}{2q - 2q\beta + 3} \right\} = \frac{3q}{2q - 2q\beta + 3} \quad (2.28)$$

and hence, in particular,

$$\lambda > \frac{3}{2 - 2\beta + \frac{3}{q}} > \frac{3}{2 - 2\beta + (2\beta - 1)} = 3 > 1,$$

and moreover has the property that

$$\lambda < \frac{3q}{2q - 2q\beta} = \frac{3}{2(1 - \beta)} \leq p, \quad (2.29)$$

so that

$$\|\varphi\|_{L^\lambda(\Omega)} \leq \|\varphi\|_{L^p(\Omega)}^{\frac{p(\lambda-1)}{(p-1)\lambda}} \|\varphi\|_{L^1(\Omega)}^{\frac{p-\lambda}{(p-1)\lambda}} \quad \text{for all } \varphi \in L^p(\Omega) \quad (2.30)$$

by the Hölder inequality.

We next observe that, again by (2.16), $\eta_0 := 2q\beta - q - 3$ is positive and thus

$$\psi(\widehat{\eta}) := \left\{ \frac{2q\beta - q - 3 + \widehat{\eta}}{3q} + \widehat{\eta} + \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta} \right\} \cdot \frac{2q\beta}{2q\beta - q - 3 - \widehat{\eta}}, \quad \widehat{\eta} \in (0, \eta_0), \quad (2.31)$$

well-defined with

$$\begin{aligned}\psi(\widehat{\eta}) &\searrow \left\{ \frac{2q\beta - q - 3}{3q} + \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta} \right\} \cdot \frac{2q\beta}{2q\beta - q - 3} \\ &= \frac{2\beta}{3} + \frac{q-1}{3q} = \frac{2q\beta + q - 1}{3q} \quad \text{as } \widehat{\eta} \searrow 0,\end{aligned}$$

whence given $\eta > 0$ we can pick $\eta_1 = \eta_1(p, q, \beta, \eta) \in (0, \eta_0)$ such that

$$\psi(\eta_1) \leq \frac{2q\beta + q - 1}{3q} + \eta, \quad (2.32)$$

where we can clearly achieve that, simultaneously,

$$\eta_1 < 2q - 2q\beta. \quad (2.33)$$

According to the latter, namely, setting

$$r = r(p, q, \beta, \eta) := \frac{3q}{2q - 2q\beta + 3 - \eta_1} \quad (2.34)$$

introduces a well-defined positive number r which satisfies

$$r < \frac{3q}{2q - 2q\beta + 3 - (2q - 2q\beta)} = q, \quad (2.35)$$

and for which we moreover have

$$r > \frac{3q}{2q - 2q\beta + 3} > 3, \quad (2.36)$$

again thanks to (2.16).

We now recall known smoothing properties of the semigroup $(e^{-tB})_{t \geq 0}$ ([15], [39]) to, firstly, fix $C_1 > 0$ such that

$$\|\nabla e^{-tB} \varphi\|_{L^\infty(\Omega)} \leq C_1 \|\varphi\|_{W^{1,\infty}(\Omega)} \quad \text{for all } \varphi \in W^{1,\infty}(\Omega), \quad (2.37)$$

to, secondly, see that since $1 < \lambda \leq q$ by (2.29) and (2.27), there exists $C_2 = C_2(p, q, \beta) > 0$ fulfilling

$$\|B^\beta e^{-tB} \varphi\|_{L^q(\Omega)} \leq C_2 \cdot (1 + t^{-\beta - \frac{3}{2}(\frac{1}{\lambda} - \frac{1}{q})}) e^{-t} \|\varphi\|_{L^\lambda(\Omega)} \quad \text{for all } \varphi \in C^0(\overline{\Omega}) \text{ and } t > 0, \quad (2.38)$$

and to, thirdly, observe that as $1 < r < q$ by (2.36) and (2.35), we can similarly find $C_3 = C_3(p, q, \beta, \eta) > 0$ such that

$$\|B^\beta e^{-tB} \varphi\|_{L^q(\Omega)} \leq C_3 \cdot (1 + t^{-\beta - \frac{3}{2}(\frac{1}{r} - \frac{1}{q})}) e^{-t} \|\varphi\|_{L^r(\Omega)} \quad \text{for all } \varphi \in C^0(\overline{\Omega}) \text{ and } t > 0. \quad (2.39)$$

Apart from that, based on the fact that (2.36) actually even warrants that $r > 3$, we may employ Corollary 2.1 and furthermore invoke Lemma 2.3 to infer that whenever $m > 0$ and $K > 0$, we can find $C_i = C_i(m, K, p, q, \beta, \eta) > 0$, $i \in \{4, 5\}$, such that given $T \in (0, \infty]$ as well as f, z_0, v_0, z, v and Q satisfying (1.9), (1.8), (1.6) and (1.7), we have

$$\|v(\cdot, t)\|_{L^r(\Omega)} \leq C_4 M_p^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta_1)}(t) \quad \text{for all } t \in (0, T) \quad (2.40)$$

and

$$\left\| \nabla \left(z(\cdot, t) - e^{-tB} z_0 \right) \right\|_{L^\infty(\Omega)} \leq C_5 M_p^{\frac{p}{p-1} \cdot \frac{(q-1)(2q\beta-q-3)}{6q^2\beta}}(t) \cdot N_{q,\beta}^{\frac{q+3+\eta_1}{2q\beta}}(t) \quad \text{for all } t \in (0, T), \quad (2.41)$$

with $(M_p(t))_{t \in (0, T)}$ and $(N_{q,\beta}(t))_{t \in (0, T)}$ as given by (2.1) and (2.24).

We henceforth let $T \in (0, \infty]$ as well as f, z_0, v_0, z, v and Q be given such that (1.9), (1.8), (1.6) and (1.7) hold with some $m > 0$ and $K > 0$, and accordingly define $(M_p(t))_{t \in (0, T)}$ and $(N_{q,\beta}(t))_{t \in (0, T)}$ through (2.1) and (2.24). Then relying on a Duhamel representation associated with the first subproblem in (1.7), we can use (2.38) and (2.39) to estimate

$$\begin{aligned} & \left\| B^\beta \left(z(\cdot, t) - e^{-tB} z_0 \right) \right\|_{L^q(\Omega)} \\ &= \left\| \int_0^t B^\beta e^{-(t-s)B} f(\cdot, s) ds - \int_0^t B^\beta e^{-(t-s)B} \left\{ v(\cdot, t) \cdot \nabla z(\cdot, s) \right\} ds \right\|_{L^q(\Omega)} \\ &\leq C_2 \int_0^t \left(1 + (t-s)^{-\beta-\frac{3}{2}(\frac{1}{\lambda}-\frac{1}{q})} \right) e^{-(t-s)} \|f(\cdot, s)\|_{L^\lambda(\Omega)} ds \\ &\quad + C_3 \int_0^t \left(1 + (t-s)^{-\beta-\frac{3}{2}(\frac{1}{r}-\frac{1}{q})} \right) e^{-(t-s)} \|v(\cdot, s) \cdot \nabla z(\cdot, s)\|_{L^r(\Omega)} ds \end{aligned} \quad (2.42)$$

for $t \in (0, T)$, where by (2.30) and (1.9),

$$\|f(\cdot, t)\|_{L^\lambda(\Omega)} \leq m^{\frac{p-\lambda}{(p-1)\lambda}} M_p^{\frac{p(\lambda-1)}{(p-1)\lambda}}(t) \quad \text{for all } t \in (0, T) \text{ and } s \in (0, t), \quad (2.43)$$

and where thanks to (2.40), (2.41), (2.37), (1.8) and the fact that $M_p \geq 1$ and $N_{q,\beta} \geq 1$, for all $t \in (0, T)$ and $s \in (0, t)$ we have

$$\begin{aligned} \|v(\cdot, s) \cdot \nabla z(\cdot, s)\|_{L^r(\Omega)} &\leq \|v(\cdot, s)\|_{L^r(\Omega)} \|\nabla z(\cdot, s)\|_{L^\infty(\Omega)} \\ &\leq \|v(\cdot, s)\|_{L^r(\Omega)} \cdot \left\{ \left\| \nabla \left(z(\cdot, s) - e^{-sB} z_0 \right) \right\|_{L^\infty(\Omega)} + \|\nabla e^{-sB} z_0\|_{L^\infty(\Omega)} \right\} \\ &\leq C_4 M_p^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta_1)}(t) \cdot \left\{ C_5 M_p^{\frac{p}{p-1} \cdot \frac{(q-1)(2q\beta-q-3)}{6q^2\beta}}(t) \cdot N_{q,\beta}^{\frac{q+3+\eta_1}{2q\beta}}(t) + C_1 K \right\} \\ &\leq C_6 M_p^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta_1 + \frac{(q-1)(2q\beta-q-3)}{6q^2\beta})}(t) \cdot N_{q,\beta}^{\frac{q+3+\eta_1}{2q\beta}}(t) \end{aligned} \quad (2.44)$$

with $C_6 = C_6(m, K, p, q, \beta, \eta) := C_4 \cdot \max\{C_5, C_1 K\}$. Since

$$C_7 = C_7(p, q, \beta) := \int_0^\infty \left(1 + \sigma^{-\beta-\frac{3}{2}(\frac{1}{\lambda}-\frac{1}{q})} \right) e^{-\sigma} d\sigma \quad \text{and} \quad C_8 = C_8(p, q, \beta, \eta) := \int_0^\infty \left(1 + \sigma^{-\beta-\frac{3}{2}(\frac{1}{r}-\frac{1}{q})} \right) e^{-\sigma} d\sigma$$

are both finite due to the circumstance that

$$\beta + \frac{3}{2} \left(\frac{1}{\lambda} - \frac{1}{q} \right) < \beta + \frac{3}{2} \cdot \left(\frac{2q - 2q\beta + 3}{3q} - \frac{1}{q} \right) = 1$$

by (2.28), and that, similarly, $\beta + \frac{3}{2}(\frac{1}{r} - \frac{1}{q}) < 1$ by the first inequality in (2.36), from (2.42)-(2.44) we thus conclude that for all $t \in (0, T)$,

$$\left\| B^\beta \left(z(\cdot, t) - e^{-tB} z_0 \right) \right\|_{L^q(\Omega)} \leq C_9 M_p^{\frac{p(\lambda-1)}{(p-1)\lambda}}(t) + C_9 M_p^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta_1 + \frac{(q-1)(2q\beta-q-3)}{6q^2\beta})}(t) \cdot N_{q,\beta}^{\frac{q+3+\eta_1}{2q\beta}}(t) \quad (2.45)$$

with

$$C_9 = C_9(m, K, p, q, \beta, \eta) := \max \left\{ C_2 C_7 m^{\frac{p-\lambda}{(p-1)\lambda}}, C_3 C_6 C_8 \right\}.$$

We may now rely on the inequality $\frac{q+3+\eta_1}{2q\beta} < \frac{q+3+(2q\beta-q-3)}{2q\beta} = 1$, as asserted by our restriction that $\eta_1 < \eta_0$, to see that due to Young's inequality there exists $C_{10} = C_{10}(m, K, p, q, \beta, \eta) > 0$ such that

$$ab \leq \frac{1}{2} a^{\frac{2q\beta}{q+3+\eta_1}} + C_{10} b^{\frac{2q\beta}{2q\beta-q-3-\eta_1}} \quad \text{for all } a \geq 0 \text{ and } b \geq 0.$$

From (2.45) we therefore obtain that for all $t \in (0, T)$,

$$\begin{aligned} N_{q,\beta}(t) &\leq 1 + C_9 M_p^{\frac{p(\lambda-1)}{(p-1)\lambda}}(t) \\ &\quad + \frac{1}{2} N_{q,\beta}(t) + C_{10} \cdot \left\{ C_9 M_p^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta_1 + \frac{(q-1)(2q\beta-q-3)}{6q^2\beta})}(t) \right\}^{\frac{2q\beta}{2q\beta-q-3-\eta_1}}, \end{aligned}$$

and that thus

$$N_{q,\beta}(t) \leq 2 + 2C_9 M_p^{\frac{p(\lambda-1)}{(p-1)\lambda}}(t) \quad (2.46)$$

$$+ 2C_9^{\frac{2q\beta}{2q\beta-q-3-\eta_1}} C_{10} M_p^{\frac{p}{p-1} \cdot (\frac{r-3}{3r} + \eta_1 + \frac{(q-1)(2q\beta-q-3)}{6q^2\beta}) \cdot \frac{2q\beta}{2q\beta-q-3-\eta_1}}(t) \quad \text{for all } t \in (0, T). \quad (2.47)$$

Since

$$\frac{\lambda-1}{\lambda} = 1 - \frac{1}{\lambda} \leq 1 - \frac{2q-2q\beta+1}{3q} = \frac{2q\beta+q-1}{3q} \leq \frac{2q\beta+q-1}{3q} + \eta$$

thanks to (2.27), and since (2.34) says that

$$\frac{r-3}{3r} = \frac{1}{3} - \frac{1}{r} = \frac{1}{3} - \frac{2q-2q\beta+3-\eta_1}{3q} = \frac{2q\beta-q-3+\eta_1}{3q},$$

and that therefore

$$\left(\frac{r-3}{3r} + \eta_1 + \frac{(q-1)(2q\beta-q-3)}{6q^2\beta} \right) \cdot \frac{2q\beta}{2q\beta-q-3-\eta_1} = \psi(\eta_1) \leq \frac{2q\beta+q-1}{3q} + \eta$$

according to (2.32), according to the definitions of $(M_p(t))_{t \in (0, T)}$ and $(N_{q,\beta}(t))_{t \in (0, T)}$ in (2.1) and (2.24) we directly infer (2.26) from (2.46). \square

Our main result on (1.7) can now be achieved by suitably adjusting the auxiliary parameters q and β in the above, and by adequately coping with the correction term $e^{-tB} z_0$ in (2.26):

PROOF of Theorem 1.2. Since $\frac{2p-3}{2p} = 1 - \frac{3}{2p} > \frac{1}{2}$ due to our assumption that $p > 3$, we can

fix $\beta = \beta(p) > \frac{1}{2}$ such that $\beta \leq \frac{2p-3}{2p}$, and that thus (2.25) holds. Given $\eta > 0$, we thereupon pick $q = q(p) > \frac{3}{2\beta-1}$ suitably large fulfilling

$$\frac{2(q+1)}{3q} < \frac{2}{3} + \eta,$$

which ensures that

$$\psi(\widehat{\eta}) := \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta} + \left(\frac{2q\beta + q - 1}{3q} + \widehat{\eta} \right) \cdot \left(\frac{q+3}{2q\beta} + \widehat{\eta} \right), \quad \widehat{\eta} > 0, \quad (2.48)$$

satisfies

$$\psi(\widehat{\eta}) \searrow \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta} + \frac{(q+3)(2q\beta + q - 1)}{6q^2\beta} = \frac{2(q+1)}{3q} < \frac{2}{3} + \eta,$$

and that hence we can choose $\eta_1 = \eta_1(p, \eta) > 0$ in such a way that

$$\psi(\eta_1) \leq \frac{2}{3} + \eta. \quad (2.49)$$

We now employ Lemma 2.3 and Lemma 2.4 to see that given $m > 0$ and $K > 0$ we can find $C_1 = C_1(m, K, p, \eta) > 0$ such that whenever $T \in (0, \infty]$ and f as well as z_0, v_0, z, v and Q comply with (1.9), (1.8), (1.6) and (1.7), as before letting $(M_p(t))_{t \in (0, T)}$ and $(N_{q, \beta}(t))_{t \in (0, T)}$ be as defined in (2.1) and (2.24) we have

$$\|z(\cdot, t) - e^{-tB} z_0\|_{W^{1, \infty}(\Omega)} \leq C_1 M_p^{\frac{p}{p-1} \cdot \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta}}(t) \cdot N_{q, \beta}^{\frac{q+3}{2q\beta} + \eta_1}(t) \quad \text{for all } t \in (0, T) \quad (2.50)$$

and

$$N_{q, \beta}(t) \leq C_2 M_p^{\frac{p}{p-1} \cdot \left(\frac{2q\beta + q - 1}{3q} + \eta_1 \right)}(t) \quad \text{for all } t \in (0, T). \quad (2.51)$$

Finally fixing $C_3 > 0$ such that, in accordance with a known boundedness feature of the Neumann heat semigroup,

$$\|e^{-tB} \varphi\|_{W^{1, \infty}(\Omega)} \leq C_3 \|\varphi\|_{W^{1, \infty}(\Omega)} \quad \text{for all } \varphi \in W^{1, \infty}(\Omega),$$

we may combine (2.50) with (2.51) and (1.8) to infer that given any such m, K, T, f, z_0, v_0, z and v , and correspondingly taking $(M_p(t))_{t \in (0, T)}$ and $(N_{q, \beta}(t))_{t \in (0, T)}$ from (2.1) and (2.24), we can estimate

$$\begin{aligned} \|z(\cdot, t)\|_{W^{1, \infty}(\Omega)} &\leq \|z(\cdot, t) - e^{-tB} z_0\|_{W^{1, \infty}(\Omega)} + \|e^{-tB} z_0\|_{W^{1, \infty}(\Omega)} \\ &\leq C_1 M_p^{\frac{p}{p-1} \cdot \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta}}(t) \cdot N_{q, \beta}^{\frac{q+3}{2q\beta} + \eta_1}(t) + C_3 \|z_0\|_{W^{1, \infty}(\Omega)} \\ &\leq C_1 C_2^{\frac{q+3}{2q\beta} + \eta_1} M_p^{\frac{p}{p-1} \cdot \left\{ \frac{(q-1)(2q\beta - q - 3)}{6q^2\beta} + \left(\frac{2q\beta + q - 1}{3q} + \eta_1 \right) \cdot \left(\frac{q+3}{2q\beta} + \eta_1 \right) \right\}}(t) + C_3 K \\ &\leq C_1 C_2^{\frac{q+3}{2q\beta} + \eta_1} M_p^{\frac{p}{p-1} \cdot \psi(\eta_1)}(t) + C_3 K \quad \text{for all } t \in (0, T) \end{aligned}$$

by (2.48). In view of (2.49), this immediately establishes (1.11). \square

3 Boundedness in (1.5) for all subcritical nonlinearities. Proof of Theorem 1.3

Next addressing the flux-limited Keller-Segel-Stokes system (1.5), let us first recall standard theory on problems of related types to state the following basic result on local existence and extensibility, actually available for fairly general F .

Lemma 3.1 *If $\Phi \in W^{2,\infty}(\Omega)$ and $F \in C^2([0, \infty))$, and if n_0, c_0 and u_0 comply with (1.13), then there exist $T_{max} \in (0, \infty]$ and unique functions*

$$\begin{cases} n \in C^0(\overline{\Omega} \times [0, T_{max})) \cap C^{2,1}(\overline{\Omega} \times (0, T_{max})), \\ c \in \bigcap_{q>3} C^0([0, T_{max}); W^{1,q}(\Omega)) \cap C^{2,1}(\overline{\Omega} \times (0, T_{max})) \quad \text{and} \\ u \in \bigcup_{\vartheta \in (\frac{3}{4}, 1)} C^0([0, T_{max}); D(A^\vartheta)) \cap C^{2,1}(\overline{\Omega} \times (0, T_{max}); \mathbb{R}^3) \end{cases}$$

with the properties that $n \geq 0$ and $c \geq 0$ in $\Omega \times (0, T_{max})$, that one can find $P \in C^{1,0}(\Omega \times (0, T_{max}))$ such that (n, c, u, P) is a classical solution of (1.5) in $\Omega \times (0, T_{max})$, and that

if $T_{max} = \infty$, then for all $\vartheta \in (\frac{3}{4}, 1)$,

$$\limsup_{t \nearrow T_{max}} \left\{ \|n(\cdot, t)\|_{L^\infty(\Omega)} + \|c(\cdot, t)\|_{W^{1,\infty}(\Omega)} + \|A^\vartheta u(\cdot, t)\|_{L^2(\Omega)} \right\} = \infty. \quad (3.1)$$

For this solution we additionally have

$$\int_{\Omega} n(\cdot, t) = \int_{\Omega} n_0 \quad \text{for all } t \in (0, T_{max}). \quad (3.2)$$

PROOF. This can be seen by standard arguments well-documented in closely related contexts (cf. [40], for instance). \square

Now the reward for all our efforts related to Theorem 1.2 consists in the circumstance that its outcome facilitates the derivation of L^p bounds for the first solution component in (1.5) through a noticeably short argument which in essence reduces to quite a straightforward combination of a standard testing procedure with the conditional estimate in (1.11).

Lemma 3.2 *Suppose that $\Phi \in W^{2,\infty}(\Omega)$, that $F \in C^2([0, \infty))$ satisfies (1.12) with some $K_F > 0$ and $\alpha > \frac{1}{2}$, and that (1.13) holds. Then the solution of (1.5) from Lemma 3.1 has the property that*

$$\sup_{t \in (0, T_{max})} \|n(\cdot, t)\|_{L^p(\Omega)} < \infty \quad \text{for all } p > 3. \quad (3.3)$$

PROOF. Without loss of generality assuming that $\alpha < 1$, we note that since $2 \cdot (1 - \alpha) < 1$ by hypothesis, we can choose $\eta > 0$ such that still

$$(2 + 3\eta) \cdot (1 - \alpha) < 1. \quad (3.4)$$

Taking $\vartheta \in (\frac{3}{4}, 1)$ such that $u_0 \in D(A^\vartheta)$, for fixed $p > 3$ we may then apply Theorem 1.2 to $m := 1 + \int_{\Omega} n_0 > 0$, $K := 1 + \max\{\|c_0\|_{W^{1,\infty}(\Omega)}, \|A^\vartheta u_0\|_{L^2(\Omega)}\} > 0$, $(f, z, v, Q) := (n, c, u, P)$ and $T := T_{max}$ to find $C_1 = C_1(p) > 0$ such that

$$\|\nabla c(\cdot, t)\|_{L^\infty(\Omega)} \leq C_1 M_p^{\frac{p}{p-1} \cdot (\frac{2}{3} + \eta)}(t) \quad \text{for all } t \in (0, T_{max}), \quad (3.5)$$

where in line with (2.1), we have set $M_p(t) := 1 + \sup_{s \in (0,t)} \|n(\cdot, s)\|_{L^p(\Omega)}$ for $t \in (0, T_{max})$.

To derive (3.3) from this, we test the first equation in (1.5) against n^{p-1} and integrate by parts to obtain that since $\nabla \cdot u = 0$, due to Young's inequality we have

$$\begin{aligned} \frac{1}{p} \frac{d}{dt} \int_{\Omega} n^p + (p-1) \int_{\Omega} n^{p-2} |\nabla n|^2 &= (p-1) \int_{\Omega} n^{p-1} F(|\nabla c|^2) \nabla n \cdot \nabla c \\ &\leq \frac{p-1}{2} \int_{\Omega} n^{p-2} |\nabla n|^2 + \frac{p-1}{2} \int_{\Omega} n^p F^2(|\nabla c|^2) |\nabla c|^2 \\ &\leq \frac{p-1}{2} \int_{\Omega} n^{p-2} |\nabla n|^2 + \frac{(p-1)K_F^2}{2} \int_{\Omega} n^p |\nabla c|^{2-2\alpha} \end{aligned}$$

for all $t \in (0, T_{max})$, so that

$$\frac{d}{dt} \int_{\Omega} n^p + \frac{2(p-1)}{p} \int_{\Omega} |\nabla n^{\frac{p}{2}}|^2 \leq \frac{p(p-1)K_F^2}{2} \int_{\Omega} n^p |\nabla c|^{2-2\alpha} \quad \text{for all } t \in (0, T_{max}). \quad (3.6)$$

Here since we are assuming $2 - 2\alpha$ to be positive, we may utilize (3.5) to estimate

$$\begin{aligned} \frac{p(p-1)K_F^2}{2} \int_{\Omega} n^p |\nabla c|^{2-2\alpha} &\leq \frac{p(p-1)K_F^2}{2} \|\nabla c\|_{L^\infty(\Omega)}^{2-2\alpha} \int_{\Omega} n^p \\ &\leq C_2 M_p^{\frac{p}{p-1} \cdot (\frac{2}{3} + \eta) \cdot (2-2\alpha)}(t) \int_{\Omega} n^p \quad \text{for all } t \in (0, T_{max}), \end{aligned} \quad (3.7)$$

where $C_2 = C_2(p) := \frac{p(p-1)K_F^2}{2} \cdot C_1^{2-2\alpha}$. Now thanks to (3.2), an interpolation on the basis of the Gagliardo-Nirenberg inequality shows that with some $C_3 = C_3(p) > 0$ and $C_4 = C_4(p) > 0$ we have

$$\begin{aligned} \left\{ \int_{\Omega} n^p \right\}^{\frac{3p-1}{3(p-1)}} &= \|n^{\frac{p}{2}}\|_{L^2(\Omega)}^{\frac{2(3p-1)}{3(p-1)}} \\ &\leq C_3 \|\nabla n^{\frac{p}{2}}\|_{L^2(\Omega)}^2 \|n^{\frac{p}{2}}\|_{L^{\frac{2}{p}}(\Omega)}^{\frac{4}{3(p-1)}} + C_3 \|n^{\frac{p}{2}}\|_{L^{\frac{2}{p}}(\Omega)}^{\frac{2(3p-1)}{3(p-1)}} \\ &\leq C_4 \|\nabla n^{\frac{p}{2}}\|_{L^2(\Omega)}^2 + C_4 \quad \text{for all } t \in (0, T_{max}), \end{aligned}$$

and that thus

$$\frac{2(p-1)}{p} \int_{\Omega} |\nabla n^{\frac{p}{2}}|^2 \geq C_5 \cdot \left\{ \int_{\Omega} n^p \right\}^{\kappa} - C_4 C_5 \quad \text{for all } t \in (0, T_{max}) \quad (3.8)$$

with $C_5 = C_5(p) := \frac{2(p-1)}{pC_4}$ and

$$\kappa = \kappa(p) := \frac{3p-1}{3(p-1)} > 1. \quad (3.9)$$

Now a combination of (3.8) with (3.7) and Young's inequality shows that for each $T \in (0, T_{max})$, (3.6) entails the inequality

$$\begin{aligned}
\frac{d}{dt} \int_{\Omega} n^p + C_5 \cdot \left\{ \int_{\Omega} n^p \right\}^{\kappa} &\leq C_2 M_p^{\frac{p}{p-1} \cdot (\frac{2}{3} + \eta) \cdot (2-2\alpha)}(T) \int_{\Omega} n^p + C_4 C_5 \\
&= \left\{ \frac{C_5}{2} \cdot \left\{ \int_{\Omega} n^p \right\}^{\kappa} \right\}^{\frac{1}{\kappa}} \cdot \left\{ \left(\frac{2}{C_5} \right)^{\frac{1}{\kappa}} \cdot C_2 M_p^{\frac{p}{p-1} \cdot (\frac{2}{3} + \eta) \cdot (2-2\alpha)}(T) \right\} + C_4 C_5 \\
&\leq \frac{C_5}{2} \cdot \left\{ \int_{\Omega} n^p \right\}^{\kappa} \\
&\quad + C_6 M_p^{\frac{\kappa}{\kappa-1} \cdot \frac{p}{p-1} \cdot (\frac{2}{3} + \eta) \cdot (2-2\alpha)}(T) + C_4 C_5 \quad \text{for all } t \in (0, T)
\end{aligned}$$

with $C_6 = C_6(p) := \left\{ \left(\frac{2}{C_5} \right)^{\frac{1}{\kappa}} C_2 \right\}^{\frac{\kappa}{\kappa-1}}$. Since $M_p \geq 1$, this implies that if we let $C_7 = C_7(p) := C_6 + C_4 C_5$, then for any such T we have

$$\frac{d}{dt} \int_{\Omega} n^p + \frac{C_5}{2} \cdot \left\{ \int_{\Omega} n^p \right\}^{\kappa} \leq C_7 M_p^{\frac{\kappa}{\kappa-1} \cdot \frac{p}{p-1} \cdot (\frac{2}{3} + \eta) \cdot (2-2\alpha)}(T) \quad \text{for all } t \in (0, T),$$

which through an ODE comparison argument guarantees that

$$\int_{\Omega} n^p(\cdot, t) \leq \max \left\{ \int_{\Omega} n_0^p, \left\{ \frac{2C_7}{C_5} M_p^{\frac{\kappa}{\kappa-1} \cdot \frac{p}{p-1} \cdot (\frac{2}{3} + \eta) \cdot (2-2\alpha)}(T) \right\}^{\frac{1}{\kappa}} \right\} \quad \text{for all } t \in (0, T),$$

and that therefore, again since $M_p \geq 1$,

$$\begin{aligned}
M_p(T) &\leq 1 + \max \left\{ \|n_0\|_{L^p(\Omega)}, \left(\frac{2C_7}{C_5} \right)^{\frac{1}{p\kappa}} M_p^{\lambda}(T) \right\} \\
&\leq C_8 M_p^{\lambda}(T) \quad \text{for all } T \in (0, T_{max})
\end{aligned} \tag{3.10}$$

with $C_8 = C_8(p) := 1 + \max \left\{ \|n_0\|_{L^p(\Omega)}, \left(\frac{2C_7}{C_5} \right)^{\frac{1}{p\kappa}} \right\}$ and $\lambda = \lambda(p) := \frac{1}{\kappa-1} \cdot \frac{1}{p-1} \cdot \left(\frac{2}{3} + \eta \right) \cdot (2-2\alpha)$.

It now only remains to observe that according to (3.9), our restriction on η in (3.4) ensures that

$$\lambda = \frac{3(p-1)}{2} \cdot \frac{1}{p-1} \cdot \left(\frac{2}{3} + \eta \right) \cdot (2-2\alpha) = (2+3\eta) \cdot (1-\alpha) < 1$$

to finally conclude from (3.10) that

$$M_p(T) \leq C_8^{\frac{1}{1-\lambda}} \quad \text{for all } T \in (0, T_{max}),$$

which implies (3.3) upon taking $T \nearrow T_{max}$. \square

In view of Theorem 1.2 and Proposition 1.1, the latter immediately implies bounds for the quantities in (3.1) related to the signal concentration and the fluid velocity.

Lemma 3.3 *If $\Phi \in W^{2,\infty}(\Omega)$ and $F \in C^2([0, \infty))$ is such that (1.12) is valid with some $K_F > 0$ and $\alpha > \frac{1}{2}$, and if (1.13) holds, then there exists $\vartheta > \frac{3}{4}$ such that T_{max} as well as the functions c and u from Lemma 3.1 satisfy*

$$\sup_{t \in (0, T_{max})} \|c(\cdot, t)\|_{W^{1,\infty}(\Omega)} < \infty \quad (3.11)$$

and

$$\sup_{t \in (0, T_{max})} \|A^\vartheta u(\cdot, t)\|_{L^2(\Omega)} < \infty. \quad (3.12)$$

PROOF. The boundedness feature in (3.11) is an immediate consequence of Theorem 1.2 when applied to any fixed $p > 3$ and combined with Lemma 3.2, while (3.12) similarly results from Proposition 1.1. \square

Now thanks to the L^∞ bounds for ∇c and u implied by Lemma 3.3 due to the continuity of the embedding $D(A^\vartheta) \hookrightarrow L^\infty(\Omega; \mathbb{R}^3)$, a straightforward application of heat semigroup estimates to the first equation in (1.5) finally yields L^∞ estimates also for n :

Lemma 3.4 *Let $\Phi \in W^{2,\infty}(\Omega)$ and $F \in C^2([0, \infty))$ be such that (1.12) holds with some $K_F > 0$ and $\alpha > \frac{1}{2}$, and assume (1.13). Then with T_{max} and n taken from Lemma 3.1, we have*

$$\sup_{t \in (0, T_{max})} \|n(\cdot, t)\|_{L^\infty(\Omega)} < \infty. \quad (3.13)$$

PROOF. We fix any $\lambda > 3$ and then readily infer from (1.12), Lemma 3.2 and Lemma 3.3 the existence of $C_1 > 0$ and $C_2 > 0$ such that $h_1 := nF(|\nabla c|^2)\nabla c + nu$ and $h_2 := n$ satisfy

$$\|h_1(\cdot, t)\|_{L^\lambda(\Omega)} \leq C_1 \quad \text{and} \quad \|h_2(\cdot, t)\|_{L^{\frac{\lambda}{2}}(\Omega)} \leq C_2 \quad \text{for all } t \in (0, T_{max}).$$

As known smoothing estimates for the Neumann heat semigroup $(e^{t\Delta})_{t \geq 0}$ on Ω ([16], [39]) provide $C_3 > 0$ and $C_4 > 0$ fulfilling

$$\|e^{t\Delta} \nabla \cdot \varphi\|_{L^\infty(\Omega)} \leq C_3 \cdot (1 + t^{-\frac{1}{2} - \frac{3}{2\lambda}}) \|\varphi\|_{L^\lambda(\Omega)} \quad \text{for all } \varphi \in C^1(\bar{\Omega}; \mathbb{R}^3) \text{ such that } \varphi \cdot \nu = 0 \text{ on } \partial\Omega$$

as well as

$$\|e^{t\Delta} \varphi\|_{L^\infty(\Omega)} \leq C_4 \cdot (1 + t^{-\frac{3}{\lambda}}) \|\varphi\|_{L^{\frac{\lambda}{2}}(\Omega)} \quad \text{for all } \varphi \in C^0(\bar{\Omega}),$$

by means of a variation-of-constants representation associated with the identity $n_t = \Delta n - n - \nabla \cdot h_1 + h_2$ we thus obtain that due to the maximum principle,

$$\begin{aligned} \|n(\cdot, t)\|_{L^\infty(\Omega)} &= \left\| e^{t(\Delta-1)} n_0 - \int_0^t e^{(t-s)(\Delta-1)} \nabla \cdot h_1(\cdot, s) ds + \int_0^t e^{(t-s)(\Delta-1)} h_2(\cdot, s) ds \right\|_{L^\infty(\Omega)} \\ &\leq e^{-t} \|n_0\|_{L^\infty(\Omega)} + C_3 \int_0^t \left(1 + (t-s)^{-\frac{1}{2} - \frac{3}{2\lambda}}\right) e^{-(t-s)} \|h_1(\cdot, s)\|_{L^\lambda(\Omega)} ds \\ &\quad + C_4 \int_0^t \left(1 + (t-s)^{-\frac{3}{\lambda}}\right) e^{-(t-s)} \|h_2(\cdot, s)\|_{L^{\frac{\lambda}{2}}(\Omega)} ds \\ &\leq \|n_0\|_{L^\infty(\Omega)} + C_1 C_3 \int_0^\infty (1 + \sigma^{-\frac{1}{2} - \frac{3}{2\lambda}}) e^{-\sigma} d\sigma \\ &\quad + C_2 C_4 \int_0^\infty (1 + \sigma^{-\frac{3}{\lambda}}) e^{-\sigma} d\sigma \quad \text{for all } t \in (0, T_{max}). \end{aligned}$$

The claim thus follows from the observation that the rightmost two integrals herein are both finite due to the fact that the inequality $\lambda > 3$ warrants that $-\frac{1}{2} - \frac{3}{2\lambda} > -1$ and $-\frac{3}{\lambda} > -1$. \square

Our main result for (1.5) has thus actually been established already:

PROOF of Theorem 1.3. Both the statement on global existence and uniqueness and the claim concerning the boundedness property in (1.14) directly result from Lemma 3.1 in conjunction with Lemma 3.3 and Lemma 3.4. \square

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